X-Ray Diffraction Study on the Local Structure of Molten ErCl₃

Yasuhiko Iwadate, Takuma Iida, Kazuko Fukushima, Junichi Mochinaga*, and Marcelle Gaune-Escard**

Department of Materials Science, Faculty of Engineering, Chiba University, Yayoi-cho 1-33, Inage-ku, Chiba 263, Japan

Z. Naturforsch. 49a, 811-814 (1994); Received July 14, 1993

The short range structure of molten $\mathrm{ErCl_3}$ at 1053 K was investigated by X-ray diffraction. The nearest neighbour distance of $\mathrm{Er^{3}}^+-\mathrm{Cl^-}$ and the coordination number of $\mathrm{Cl^-}$ around $\mathrm{Er^{3}}^+$ were estimated to be 2.63 Å and 5.8 Å, respectively, and the $\mathrm{Er^{3}}^+-\mathrm{Er^{3}}^+$ and $\mathrm{Cl^--Cl^-}$ distances 4.05 Å and 3.75 Å, respectively. The ratio of the anion-anion pair distance to the cation-anion pair distance (r_{--}/r_{+-}) was, therefore, 1.43, being close to 1.41, characteristic of octahedral geometry. These results indicate that $\mathrm{ErCl_3^6}^-$ octahedra exist in molten $\mathrm{ErCl_3}$. Probably $\mathrm{Er_2Cl_{10}^{4-}}$ and a small amount of $\mathrm{Er_2Cl_{11}^{4-}}$ ions are formed according to the geometrical calculations of bond lengths and bond angles.

1. Introduction

Pioneering high temperature Raman spectroscopic studies have been carried out in the 1970's on rare earth chloride systems in the liquid and solid state [1-3]. Structural analyses have been made by X-ray diffraction and Raman spectroscopy on pure lanthanide trichlorides [4, 5] with UCl₃-type crystal structure (hexagonal). Those results indicated that the nearest neighbour chloride coordination number of a lanthanide ion changed from 9 in the crystal to 6 in the melt, allowing for octahedral geometry around the lanthanide ions. Little is known about melts of lanthanide trichlorides with AlCl₃-type crystal structure (monoclinic), except for YCl₃ [6]. The melting behaviour of ErCl₃ might be similar to that of YCl₃, having the same crystal structure. Recently, structural correlations in molten trivalent metal chlorides have been estimated as functions of the metal ion size for LaCl₃ and AlCl₃-type lanthanide trichlorides, using a charged soft-sphere model and the hypernetted chain approximation [7]. Furthermore, Erbölükbas et al. [8] have analyzed the available data on bond lengths and Raman frequencies by treating an isolated octahedral complex ion within a model which adopts charged soft-sphere interionic potentials, supplemented by taking ionic polarization into account.

* Formerly at Chiba University.

Reprint requests to Prof. Y. Iwadate.

In the present work, the short range structure of molten ErCl₃ was studied by X-ray diffraction. Possible anionic pairs are discussed in view of interatomic distances and bond angles.

2. Experimental

ErCl₃ was prepared by heating a mixture of analytical grade Er₂O₃ and 2.5 times the theoretical amount of NH₄Cl at 623 K for 3 h. Excess NH₄Cl was removed by heating at 1173 K. The chlorination reaction is

$$Er_2O_3 + 6 NH_4Cl \rightarrow 2 ErCl_3 + 6 NH_3 + 3 H_2O$$
.

ErCl₃ thus obtained was purified by ErCl₃ vapour deposition onto the surface of a quartz tube cooled with water under reduced pressure down to 10⁻² Pa for 12 h. The vapour source was heated at 1253 K, so as to remove small amounts of water, oxychloride, unreacted NH₄Cl and oxide. The molar volumes were measured dilatometrically using a transparent thin goldfilm-coated furnace with a good thermal stability [9]. In the X-ray scattering measurements, the purified sample was sealed in a thin quartz tube cell.

The X-ray diffractometer, having a $\theta-\theta$ type reflection geometry (Rigaku Corporation, Tokyo, Japan), was employed with MoK α radiation, and the diffracted beam was monochromatized with curved graphite. The Q-range, $0.93 \le Q/\text{Å}^{-1} \le 12.5$, was explored with two pairs of slits $[1/2^\circ-1/2^\circ]$ when $2.5^\circ \le \theta \le 15^\circ$, and $[1^\circ-1^\circ]$ when $13^\circ \le \theta \le 45^\circ$ ($2\theta = \text{scattering}$ angle and $Q = 4\pi \sin \theta/\lambda$ with λ being the

0932-0784 / 94 / 0700-0811 \$ 06.00 © - Verlag der Zeitschrift für Naturforschung, D-72027 Tübingen



Dieses Werk wurde im Jahr 2013 vom Verlag Zeitschrift für Naturforschung in Zusammenarbeit mit der Max-Planck-Gesellschaft zur Förderung der Wissenschaften e.V. digitalisiert und unter folgender Lizenz veröffentlicht: Creative Commons Namensnennung-Keine Bearbeitung 3.0 Deutschland

This work has been digitalized and published in 2013 by Verlag Zeitschrift für Naturforschung in cooperation with the Max Planck Society for the Advancement of Science under a Creative Commons Attribution-NoDerivs 3.0 Germany License.

^{**} Institut Universitaire des Systèmes Thermiques Industriels UA 1168, Université de Provence, Centre St. Jérôme, 13397 Marseille Cedex 20, France.

wavelength of the radiation). The scattered X-ray intensities so obtained were corrected for background, polarization, absorption and Compton scattering, and normalized to the coherent scattering intensity of the stoichiometric unit, in accordance with the usual procedure [10]. Atomic scattering factors and Compton scattering factors were taken from the literature [11, 12]. Thus experimental structure factors and radial distribution functions were obtained.

3. Results and Discussion

The radial distribution function D(r), the pair correlation function G(r) and the reduced intensity function $Q \cdot i(Q)$ are defined as follows:

$$Q \cdot i(Q) = Q \cdot \left[I_{\text{eu}}^{\text{coh}}(Q) - \sum_{i} f_{i}(Q)^{2} \right] / \sum_{i} f_{i}(Q)^{2}, \quad (1)$$

$$D(r) = 4 \pi r^{2} \left(\sum_{i} \overline{K}_{i} \right)^{2} \varrho_{0} \qquad (2)$$

$$+ \sum_{i} (\overline{K}_{i})^{2} (2r/\pi) \int_{0}^{Q_{\text{max}}} Q \cdot i(Q) \sin(Qr) \, dQ,$$

$$G(r) = 1 + \left[\sum_{i} (\overline{K}_{i})^{2} / \left\{ 2\pi^{2} r \left(\sum_{i} \overline{K}_{i} \right)^{2} \varrho_{0} \right\} \right]$$

 $\int_{0}^{Q_{\max}} Q \cdot i(Q) \sin(Qr) dQ,$

(3)

where ϱ_0 is the number of stoichiometric units per unit volume, \bar{K}_i the effective electron number in the atom i [13], $f_i(Q)$ the independent atomic scattering amplitude, $I_{\rm cu}^{\rm coh}(Q)$ the total coherent intensity function and $Q_{\rm max}$ the maximum value of Q reached in the experiments. Summation for i is made over the stoichiometric unit. These functions were calculated with the parameters listed in Table 1; details of the analytical procedure are described in [14].

3.1. Molar Volume and X-ray Diffraction Analyses

 $ErCl_3$ forms a monoclinic crystal [15]. In order to study the volume change on melting, the molar volume, V, in the crystalline state was calculated by the equation

$$V = abc \sin \beta (N_A/Z),$$

where a, b, c, and β are the lattice constants. N_A and Z are Avogadro's number and the number of stoichiometric units contained in the unit cell, respectively.

Table 1. Numerical parameters in the radial disribution analysis.

System	ErCl ₃
Temperature/K Density/g cm ⁻³	1053
Density/g cm ⁻³	3.619
Number density/Å ⁻³	0.007965
Effective electron number	
$ar{K}_{ ext{Er}}$	74.0714
\bar{K}_{Cl}	14.9762
$Q_{\text{max}}/\mathring{\mathrm{A}}^{-1}$	12.5

Table 2. Molar volume changes on melting.

System	$T_{\rm mp}/{ m K}$	$V_{\rm s}/{\rm cm}^3~{\rm mol}^{-1}$	$V_{\rm m}/{\rm cm}^3~{\rm mol}^{-1}$	$100 (V_{\rm m} - V_{\rm s}) V_{\rm s}^{-1} / \%$
LaCl ₃	1150	63.90	76.30	19.1
PrCl ₃	1059	61.40	74.30	21.0
NdCl ₃	1029	60.47	73.65	21.8
SmCl ₃	935	59.23	73.73	24.5
GdCl ₃	875	58.03	73.34	26.4
DyCl ₃	928	74.38	74.62	0.3
HoCl ₃	993	73.02	74.00	1.3
ErCl ₃	1046	72.11	75.65	4.9
YCl ₃	987	74.83	75.17	0.5

 $T_{\rm mp}$, $V_{\rm s}$, and $V_{\rm m}$ refer to the melting point, the molar volumes of solid and melt, respectively.

For the lanthanide trichlorides with hexagonal structure [16] the molar volume was calculated by

$$V = (\sqrt{3}/2) a^2 c(N_A/Z)$$
.

The molar volume of molten ${\rm ErCl_3}$ was measured and expressed as a linear function of temperature as $V_{\rm m} = 59.63 + 1.562 \times 10^{-2}~T$ with $V_{\rm m}$ and T in ${\rm cm^3\,mol^{-1}}$ and K, respectively. The molar volumes of several lanthanide trichlorides are given in Table 2.

As can be seen from Table 2, the changes on melting of the molar volume of UCl₃-type lanthanide trichlorides are very large, those of AlCl₃-type lanthanide trichlorides being small. Therefore, the nearest neighbour ordering in molten ErCl₃ appears to be similar to that in the solid state.

The G(r) and D(r) curves of molten $ErCl_3$ are shown in Figs. 1 and 2, respectively.

The center of the first peak in G(r) appeared in the range of 2.60 Å to 2.65 Å. Considering the sum of individual ionic radii given by Shannon [17] with Er³⁺ and Cl⁻ being 0.89 Å and 1.81 Å, respectively, the first peak was assigned to the nearest neighbour correlation, Er³⁺–Cl⁻; the other correlations hardly contributed to this peak. The coordination number of Cl⁻ ion around Er³⁺ ion was estimated at 5.8 ± 0.2 by inte-

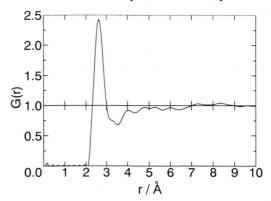


Fig. 1. Correlation function G(r).

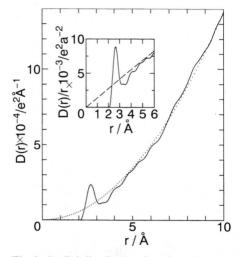


Fig. 2. Radial distribution function D(r).

gration of the area under the peak of D(r) shown in Figure 2. The second peak was found at 3.6-4.3 Å, being assignable to the Cl^--Cl^- and $Er^{3+}-Er^{3+}$ correlations. For example, the $Er^{3+}-Er^{3+}$ distance in the solid $ErCl_3$ was estimated at 3.91 Å from crystallographical data [15] (see Table 1). Some model structure was taken into account, and the reduced intensity function, $Q \cdot i(Q)$, of the model was reversely computed from the equation

$$Q \cdot i(Q) = \left[\sum_{i} \sum_{j} n_{ij} f_i(Q) f_j(Q) \exp(-b_{ij} Q^2) \right] \cdot \sin(Q r_{ij}) / r_{ij} / \sum_{i} f_i(Q)^2, \quad (4)$$

where n_{ij} , r_{ij} , and b_{ij} are the coordination number, the interatomic distance, and the temperature factor for

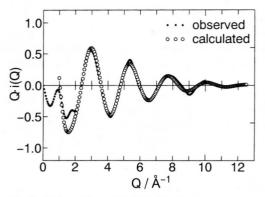


Fig. 3. Reduced intensity function $Q \cdot i(Q)$.

Table 3. Coordination number n_{ij} , interionic distance r_{ij} in crystal (averaged) [15] and melt, root mean square displacement $\langle \Delta r_{ij}^2 \rangle^{1/2}$, least squares fitted for molten ErCl₃, and temperature factor b_{ij} in (4).

i	j	n_{ij}	$r_{ij}/ ext{\AA}$		$\langle \Delta r_{ij}^2 \rangle^{1/2} / c$	$ m \mathring{A} \ \it b_{ij}/\mathring{A}^{2}$
			Crystal	Melt		
Er ³⁺	Cl-	5.8	2.58	2.63	0.23	0.03
Cl ⁻ Er ³⁺	Cl ⁻ Er ³⁺	$\begin{array}{c} 8-9 \\ 1-2 \end{array}$	3.36 3.91	3.75 4.05	0.72 0.89	0.26 0.39

the i-j pair, respectively. The structural parameters obtained from G(r) and D(r) were used as the initial values of this process for their further refinement. Each parameter was optimized by a least squares calculation beyond $Q = 1.0 \text{ Å}^{-1}$. The parameters of the most probable model in the $Q \cdot i(Q)$ fitting are given in Table 3. The observed and calculated $Q \cdot i(Q)$'s are shown in Figure 3. The R factor for the best fitted model, $R = \sum |Q \cdot i(Q)_{\text{obs}} - Q \cdot i(Q)_{\text{calc}}|/\sum |Q \cdot i(Q)_{\text{obs}}|$ at intervals of 0.05 Å⁻¹ in the range of $1.0 \le Q/\text{Å}^{-1}$ \leq 12.5, converged to 0.24. From these results, the nearest neighbour distance was estimated at 2.63 Å and the coordination number of Cl⁻ around Er³⁺ was 5.8. Like-ion pair distances of Cl⁻-Cl⁻ and Er³⁺-Er³⁺ were calculated to be 3.75 Å and 4.05 Å, respectively. The ratio of r_{--}/r_{+-} was 1.43, being close to the value for octahedral geometry 1.41.

3.2. Medium Range Structure

The interatomic structure, especially in the Er³⁺-Er³⁺ correlation, was considered in terms of molar volume and obtained short range structural parameters. If Er³⁺ in ErCl₃ melt were distributed isotropically, the

average Er³⁺-Er³⁺ pair-distance would be 5.01 Å. The corresponding distance obtained from the X-ray diffraction analyses was 4.05 Å. This indicates that, as to the Er³⁺ - Er³⁺ pair, there is some specific configuration in the melt.

In order to discuss the medium range structure of the melt, the following models were taken into account: (i) edge-sharing octahedra, (ii) corner-sharing octahedra. When undistorted octahedral anions form an edge-sharing configuration, the Er³⁺-Er³⁺ pairdistance would be about 3.73 Å at most. This value is, however, smaller than that obtained by X-ray diffraction or molar volume analyses. The rigid octahedra edge-sharing model is thus not suitable for the medium range structure of molten ErCl3, but a model assuming flexible octahedra with variable Er3+-Cl- and Cl⁻-Cl⁻ bond lengths could be valid. This supports the suggestion by Saboungi et al. [6] that molten ErCl₃ may have a structure similar to YCl₃ (see Table 3). Next, a linear arrangement of octahedral ions was considered on the basis of corner-sharing. In this model, the Er³⁺ – Er³⁺ pair-distance became 5.28 Å, a too large value. If the arrangement is not linear, i.e. $\angle Er^{3+} - Cl^{-} - Er^{3+} \approx 100^{\circ}$, and two octahedra were

twisted with each other, this structure might exist. But corner-sharing is less probable than edge-sharing because of the steric hindrance between Cl ions, although corner-sharing has been expected from numerical calculations on Raman frequencies [8].

4. Conclusions

By X-ray diffraction, the existence of octahedral complex ions, ErCl₆³⁻, in molten ErCl₃ is confirmed. Also expected is some clustering of distorted octahedra with edge-sharing.

Acknowledgements

The present work was financially supported by a research contribution from KATO Science Advancement Foundation, Université de Provence, CNRS and Conseil Régional de Provence-Alpes-Côte d'Azur. The authors wish to thank Mitsui Mining Smelting Co., Ltd. and Nippon Yttrium Co., Ltd. for supplying high purity Er₂O₃.

- [1] V. A. Maroni, E. J. Hathaway, and G. N. Papatheodorou, J. Phys. Chem. 78, 1134 (1974).
- [2] G. N. Papatheodorou, Inorg. Nucl. Chem. Lett. 11, 483 (1975).
- [3] G. N. Papatheodorou, J. Chem. Phys. 66, 2893 (1977).
- [4] J. Mochinaga, Y. Iwadate, and K. Fukushima, Mater. Sci. Forum 73-75, 147 (1991).
- [5] J. Mochinaga and K. Igarashi, Rare Earths, No. 15, 13 (1989).
- [6] M.-L. Saboungi, D. L. Price, C. Scamehorn, and M. P.
- Tosi, Europhys. Lett. 15, 283 (1991).
 [7] H. Tatlipinar, Z. Akdeniz, G. Pastore, and M. P. Tosi, J. Phys. Condens. Matter 4, 8933 (1992).
- [8] A. Erbölükbas, Z. Akdeniz and M. P. Tosi, Il Nuovo Cimento 14, 87 (1992).
- [9] J. Mochinaga, K. Igarashi, and Y. Iwadate, J. Chem. Eng. Data 30, 274 (1985).

- [10] H. Ohno, K. Furukawa, K. Igarashi, and J. Mochinaga, J. Chem. Soc. Faraday Trans. 1, 78, 1555 (1982).
- [11] International Tables for X-ray Crystallography (J. A. Ibers and W. C. Hamilton, eds.), Kynoch, Birmingham 1974, Vol. 4, p. 99.
- [12] F. Hajdu, Acta Cryst. Sec. A 27, 73 (1971).
- [13] B. E. Warren, H. Krutter, and O. Morningstar, J. Amer. Chem. Soc. 19, 202 (1936).
- [14] Y. Iwadate, K. Igarashi, T. Hattori, S. Nishiyama, K. Fukushima, J. Mochinaga, N. Igawa, and H. Ohno, J. Chem. Phys. 99, 6890 (1993).
- [15] R. W. G. Wyckoff, Crystal Structures, Vol. 2, 2nd ed., Interscience Publ., New York 1964, pp. 55-57.
- [16] ibid., pp. 77–78.
- [17] R. D. Shannon, Acta Cryst. A 32, 751 (1976).